

## Traces and Determinants

Let  $A$  be an  $n \times n$  matrix with complex entries:

$$A = \begin{pmatrix} a_{11} & a_{12} & \cdots & a_{1n} \\ a_{21} & a_{22} & \cdots & a_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1} & a_{n2} & \cdots & a_{nn} \end{pmatrix}.$$

Then

$$\operatorname{tr} A = \sum_{i=1}^n a_{ii}$$

and

$$\det A = \sum_{\sigma \in S_n} (\operatorname{sign} \sigma) a_{1\sigma(1)} a_{2\sigma(2)} \cdots a_{n\sigma(n)}.$$

Properties of trace: For  $A$  and  $B$  in  $M(n, \mathbb{C})$  and  $S$  in  $GL(n, \mathbb{C})$ ,

- $\operatorname{tr}(A + B) = \operatorname{tr} A + \operatorname{tr} B$ ;
- $\operatorname{tr}(AB) = \operatorname{tr}(BA)$ ;
- $\operatorname{tr}(SAS^{-1}) = \operatorname{tr} A$ ;
- The trace of  $A$  is the sum of the eigenvalues of  $A$ .

Properties of determinant:

- $\det(AB) = \det(BA) = (\det A)(\det B)$ ;
- $\det(SAS^{-1}) = \det A$ ;
- The determinant of  $A$  is the product of the eigenvalues of  $A$ .

Define the *exponential* of  $A$  as

$$\exp A := \sum_{k=0}^{\infty} \frac{1}{k!} A^k.$$

Warning: In general  $\exp(A + B) \neq (\exp A)(\exp B)$  unless  $A$  and  $B$  commute.

**Theorem:**  $\det(\exp A) = e^{\text{tr} A}$

**Proof:** First consider the case of a  $k \times k$  Jordan block:

$$\exp \text{tr} \begin{pmatrix} c & 1 & 0 & \cdots & 0 & 0 \\ 0 & c & 1 & \cdots & 0 & 0 \\ 0 & 0 & c & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & c & 1 \\ 0 & 0 & 0 & \cdots & 0 & c \end{pmatrix} = e^{ck}$$

$$\det \exp \begin{pmatrix} c & 1 & 0 & \cdots & 0 & 0 \\ 0 & c & 1 & \cdots & 0 & 0 \\ 0 & 0 & c & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & c & 1 \\ 0 & 0 & 0 & \cdots & 0 & c \end{pmatrix} = \det \begin{pmatrix} e^c & e & 0 & \cdots & 0 & 0 \\ 0 & e^c & e & \cdots & 0 & 0 \\ 0 & 0 & e^c & \cdots & 0 & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots & \vdots \\ 0 & 0 & 0 & \cdots & e^c & e \\ 0 & 0 & 0 & \cdots & 0 & e^c \end{pmatrix} = (e^c)^k = e^{ck}.$$

Therefore the theorem is true for Jordan blocks. Next, suppose we have a matrix in Jordan canonical form:

$$J = \begin{pmatrix} J_1 & 0 & 0 & \cdots & 0 \\ 0 & J_2 & 0 & \cdots & 0 \\ 0 & 0 & J_3 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & J_m \end{pmatrix}.$$

Then

$$\exp J = \begin{pmatrix} \exp J_1 & 0 & 0 & \cdots & 0 \\ 0 & \exp J_2 & 0 & \cdots & 0 \\ 0 & 0 & \exp J_3 & \cdots & 0 \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & 0 & \cdots & \exp J_m \end{pmatrix}$$

and so

$$\begin{aligned} \det \exp J &= (\det \exp J_1)(\det \exp J_2) \cdots (\det \exp J_n) \\ &= e^{\text{tr} J_1} e^{\text{tr} J_2} \cdots e^{\text{tr} J_n} \\ &= e^{\text{tr} J_1 + \text{tr} J_2 + \cdots + \text{tr} J_n} \\ &= e^{\text{tr} J}. \end{aligned}$$

Finally, given  $A$  in  $M(n, \mathbb{C})$ , write  $A = SJS^{-1}$ , where  $J$  is in Jordan canonical form. Then

$$\begin{aligned} \det \exp A &= \det \exp(SJS^{-1}) = \det (S(\exp J)S^{-1}) \\ &= \det \exp J = e^{\operatorname{tr} J} = e^{\operatorname{tr}(SJS^{-1})} = e^{\operatorname{tr} A}. \quad \square \end{aligned}$$

Let  $V$  be a complex vector space equipped with an *inner product*. This is a function  $\langle \cdot, \cdot \rangle : V \times V \rightarrow \mathbb{C}$  such that for all elements  $v, w$ , and  $u$  in  $V$  and all complex numbers  $\alpha$  and  $\beta$ ,

- $\langle \alpha v + \beta w, u \rangle = \alpha \langle v, u \rangle + \beta \langle w, u \rangle$ ;
- $\langle v, \alpha w + \beta u \rangle = \bar{\alpha} \langle v, w \rangle + \bar{\beta} \langle v, u \rangle$ ;
- $\langle w, v \rangle = \overline{\langle v, w \rangle}$ ;
- $\langle v, v \rangle \geq 0$ , with  $\langle v, v \rangle = 0$  if and only if  $v = 0$ .

An *orthonormal basis* for  $V$  is a vector space basis  $\{e_k\}_{k=1}^n$  for  $V$  with the additional properties

- $\langle e_k, e_k \rangle = 1$  for  $1 \leq k \leq n$ ;
- $\langle e_k, e_\ell \rangle = 0$  for  $k \neq \ell$ .

Let  $A$  be a linear transformation of  $V$ . Then

$$\operatorname{tr} A = \sum_{k=1}^n \langle Ae_k, e_k \rangle$$

and

$$\det A = \sum_{\sigma \in S_n} (\operatorname{sign} \sigma) \langle Ae_1, e_{\sigma(1)} \rangle \langle Ae_2, e_{\sigma(2)} \rangle \cdots \langle Ae_n, e_{\sigma(n)} \rangle.$$

These quantities are independent of the choice of orthonormal basis.

The *adjoint* of  $A$  is the linear transformation determined by the equation

$$\langle Av, w \rangle = \langle v, A^*w \rangle$$

for all  $v$  and  $w$  in  $V$ .

If we write  $A$  as a matrix with respect to an orthonormal basis, then  $A^*$  is the complex conjugate transpose of  $A$ ; i.e., the  $(i, j)$  entry of  $A^*$  is  $\overline{a_{ji}}$ . Thus

$$\operatorname{tr} A^* = \overline{\operatorname{tr} A}, \quad \det A^* = \overline{\det A}.$$

Now let  $V$  be an infinite-dimensional complex inner product space and define a norm  $\|v\| := \sqrt{\langle v, v \rangle}$  for every  $v$  in  $V$ . We say that  $V$  is *complete* if every Cauchy sequence with respect to this norm is convergent. In this case we will use the letter  $\mathcal{H}$  to denote our complex inner product space, and we call it a *Hilbert space*.

We will only consider *separable* Hilbert spaces. This means that  $\mathcal{H}$  contains a countably infinite subset  $\{e_k\}$  with the following properties:

- $\langle e_k, e_k \rangle = 1$  for all  $k$ ;
- $\langle e_k, e_\ell \rangle = 0$  for  $k \neq \ell$ ;
- $v = \sum_{k=1}^{\infty} \langle v, e_k \rangle e_k$  for every  $v$  in  $\mathcal{H}$ .

Warning: the set  $\{e_k\}$  is **not** a vector space basis!

Let  $A$  be a linear transformation of  $\mathcal{H}$ . We say that  $A$  is *bounded* if

$$\|A\| := \sup \left\{ \frac{\|Av\|}{\|v\|} : v \neq 0 \right\} < \infty.$$

We will call a bounded linear transformation of  $\mathcal{H}$  an *operator* on  $\mathcal{H}$ .

The collection of all operators on  $\mathcal{H}$  is an *algebra* (closed under addition, multiplication [composition], scalar multiplication), and is denoted  $\mathcal{B}(\mathcal{H})$ .

How do we define trace for operators on  $\mathcal{H}$ ?

Naive idea: choose an orthonormal basis  $\{e_k\}$  for  $\mathcal{H}$  and set

$$\text{tr } A = \sum_{k=1}^{\infty} \langle Ae_k, e_k \rangle.$$

Problem 1: The right-hand side does not necessarily converge.

Example:

$$\text{tr } I = \sum_{k=1}^{\infty} \langle Ie_k, e_k \rangle = \sum_{k=1}^{\infty} \langle e_k, e_k \rangle = \sum_{k=1}^{\infty} 1 = \infty.$$

So not every operator has a well-defined trace.

Problem 2: Even if the right-hand side does converge, its value may depend on the choice of orthonormal basis.

An operator  $P$  on  $\mathcal{H}$  is *positive* if  $\langle Pv, v \rangle \geq 0$  for all  $v$  in  $\mathcal{H}$ .

Example: Let  $A$  be any operator on  $\mathcal{H}$ . Then  $A^*A$  is positive, because

$$\langle A^*Av, v \rangle = \langle Av, Av \rangle \geq 0.$$

In fact, every positive operator  $P$  has this form for some operator  $A$ .

If  $P$  is positive, then  $\sum_{k=1}^{\infty} \langle Pe_k, e_k \rangle$  is in  $[0, \infty]$  and is independent of the choice of orthonormal basis.

Every positive operator  $P$  has a positive *square root operator*  $\sqrt{P}$ . Define

$$|A| := \sqrt{A^*A}.$$

Example: Take

$$A = \begin{pmatrix} -\frac{27}{25} + \frac{32}{25}i & -\frac{36}{25} - \frac{24}{25}i \\ -\frac{36}{25} - \frac{24}{25}i & -\frac{48}{25} + \frac{18}{25}i \end{pmatrix}.$$

Then

$$A^*A = \begin{pmatrix} \frac{29}{5} & \frac{12}{5} \\ \frac{12}{5} & \frac{36}{5} \end{pmatrix}.$$

Let

$$S = \begin{pmatrix} \frac{3}{5} & -\frac{4}{5} \\ \frac{4}{5} & \frac{3}{5} \end{pmatrix}.$$

Then

$$S^{-1}(A^*A)S = \begin{pmatrix} 9 & 0 \\ 0 & 4 \end{pmatrix},$$

whence

$$\sqrt{S^{-1}(A^*A)S} = \begin{pmatrix} 3 & 0 \\ 0 & 2 \end{pmatrix}$$

and thus

$$|A| = S \left( \sqrt{S^{-1}(A^*A)S} \right) S^{-1} = \begin{pmatrix} \frac{59}{25} & \frac{12}{25} \\ \frac{12}{25} & \frac{66}{25} \end{pmatrix}.$$

Define

$$\mathcal{L}^1(\mathcal{H}) := \left\{ A \in \mathcal{B}(\mathcal{H}) : \sum_{k=1}^{\infty} \langle |A|e_k, e_k \rangle < \infty \right\}.$$

The set  $\mathcal{L}^1(\mathcal{H})$  is an ideal in  $\mathcal{B}(\mathcal{H})$  and is called the *ideal of trace-class operators* on  $\mathcal{H}$ . For  $A$  in  $\mathcal{L}^1(\mathcal{H})$  we can define  $\text{tr } A$  in the naive way we originally proposed:

$$\text{tr } A = \sum_{k=1}^{\infty} \langle Ae_k, e_k \rangle.$$

Properties of  $\text{tr}$ :

- $\text{tr}(A + B) = \text{tr } A + \text{tr } B$  for  $A$  and  $B$  in  $\mathcal{L}^1(\mathcal{H})$ ;
- $\text{tr}(AB) = \text{tr}(BA)$  for  $A$  in  $\mathcal{L}^1(\mathcal{H})$  and  $B$  in  $\mathcal{B}(\mathcal{H})$ ;
- $\text{tr}(SAS^{-1}) = \text{tr } A$  for  $A$  in  $\mathcal{L}^1(\mathcal{H})$  and  $S$  in  $\mathcal{B}(\mathcal{H})$  invertible;
- $\text{tr } A$  is the sum of the eigenvalues of  $A$  for all  $A$  in  $\mathcal{L}^1(\mathcal{H})$ .

Remark: This last statement, known as Lidskii's theorem, was not proved until 1959.

How do we define the determinant?

For  $\|A\| < 1$ , we can define the logarithm of  $I + A$  by the infinite series

$$\log(I + A) = \sum_{n=1}^{\infty} \frac{(-1)^{n+1}}{n} A^n.$$

If  $A$  is trace class, then for  $\mu \in \mathbb{C}$  with sufficiently small modulus, the operator  $\log(I + \mu A)$  is also trace class, so we can define

$$\det(I + \mu A) = e^{\text{tr}(\log(I + \mu A))}$$

and then extend by analytic continuation, so that the domain of  $\det$  is

$$\text{GL}(1, (I + \mathcal{L}^1(\mathcal{H}))),$$

the multiplicative group of invertible elements of  $\mathcal{B}(\mathcal{H})$  of the form  $I + L$  for some  $L$  in  $\mathcal{L}^1(\mathcal{H})$ .

Properties of  $\det$ :

- $\det(AB) = (\det A)(\det B)$  for  $A$  and  $B$  in  $\text{GL}(1, I + \mathcal{L}^1(\mathcal{H}))$ ;
- $\det A^{-1} = (\det A)^{-1}$  for  $A$  in  $\text{GL}(1, (I + \mathcal{L}^1(\mathcal{H})))$ ;
- $\det(SAS^{-1}) = \det A$  for  $A$  in  $\text{GL}(1, (I + \mathcal{L}^1(\mathcal{H})))$  and  $S$  in  $\mathcal{B}(\mathcal{H})$  invertible;
- $\det A$  is the product of the eigenvalues of  $A$  for  $A$  in  $\text{GL}(1, I + \mathcal{L}^1(\mathcal{H}))$ .

These quantities are hard to compute directly, especially the determinant! However, in certain cases of geometric and/or topological interest, there are other ways to proceed.

**Example 1:**

Suppose  $K : [a, b] \times [a, b] \rightarrow \mathbb{C}$  is continuous and define  $A$  in  $\mathcal{B}(L^2[a, b])$  by the formula

$$(Af)(x) = \int_a^b K(x, y)f(y) dy.$$

This is an example of a *compact* operator. It is not always trace class (in fact, it is an open problem to find necessary and sufficient conditions on  $K$  so that  $A$  is trace class), but if  $A$  is trace class, then

$$\text{tr } A = \int_a^b K(x, x) dx.$$

We can also express  $\det(I + A)$  in terms of  $K$ . For each  $n$ -tuple  $(x_1, x_2, \dots, x_n)$  in  $[a, b]$ , define

$$K_n(x_1, x_2, \dots, x_n) = \det \begin{pmatrix} K(x_1, x_1) & K(x_1, x_2) & \cdots & K(x_1, x_n) \\ K(x_2, x_1) & K(x_2, x_2) & \cdots & K(x_2, x_n) \\ \vdots & \vdots & \ddots & \vdots \\ K(x_n, x_1) & K(x_n, x_2) & \cdots & K(x_n, x_n) \end{pmatrix}$$

Then

$$\det(I + A) = 1 + \sum_{n=1}^{\infty} \frac{1}{n!} \int_a^b \cdots \int_a^b K_n(x_1, x_2, \dots, x_n) dx_1 dx_2 \dots dx_n.$$

**Example 2:**

Consider the Hilbert space  $L^2(S^1)$  with the inner product

$$\langle f, g \rangle = \frac{1}{2\pi} \int_0^{2\pi} f(\theta) \overline{g(\theta)} d\theta.$$

This Hilbert space has orthonormal basis

$$\{e^{in\theta} : n \in \mathbb{Z}\} = \{z^n : n \in \mathbb{Z}\}.$$

Let  $C(S^1)$  denote the algebra of continuous complex-valued functions on the circle. For each  $\phi$  in  $C(S^1)$ , define an operator  $M_\phi$  on  $L^2(S^1)$  via pointwise multiplication:

$$(M_\phi f)(x) = \phi(x)f(x).$$

Next, let  $H^2(S^1)$  be the Hilbert subspace of  $L^2(S^1)$  whose orthonormal basis is

$$\{z^n : n \geq 0\}.$$

An alternate description of  $H^2(S^1)$  is the Hilbert subspace of the elements of  $L^2(S^1)$  that extend to analytic functions on the disk  $\{z \in \mathbb{C} : |z| < 1\}$ .

Define the *orthogonal projection*  $P : L^2(S^1) \rightarrow H^2(S^1)$  by

$$P \left( \sum_{n=-\infty}^{\infty} a_n z^n \right) = \sum_{n=0}^{\infty} a_n z^n.$$

Then for each  $\phi$  in  $C(T)$ , define the *Toeplitz operator*  $T_\phi$  on  $H^2(S^1)$  by the formula

$$T_\phi = PM_\phi.$$

Properties of Toeplitz operators: For  $\phi$  and  $\psi$  in  $C(S^1)$  and  $\lambda$  in  $\mathbb{C}$ ,

- $T_{\phi+\psi} = T_\phi + T_\psi$ ;
- $T_{\lambda\phi} = \lambda T_\phi$ ;
- $T_\phi^* = T_{\bar{\phi}}$ .

$T_{\phi\psi} \neq T_\phi T_\psi$  in general, but for  $\phi$  and  $\psi$  in  $C^\infty(S^1)$ , we have

$$T_\phi T_\psi - T_\psi T_\phi \in \mathcal{L}^1(\mathcal{H}).$$

Surprisingly (at first), the trace of this quantity can be nonzero. This is because  $T_\phi T_\psi$  and  $T_\psi T_\phi$  are typically not trace class operators, even though their difference is.



Example:

$$T_{z^{-3}}T_{z^3}(z^n) = z^n \text{ for all } n \geq 0$$

$$T_{z^3}T_{z^{-3}}(z^n) = \begin{cases} 0 & 0 \leq n < 3 \\ z^n & n \geq 3 \end{cases}$$

Therefore

$$\text{tr}(T_{z^{-3}}T_{z^3} - T_{z^3}T_{z^{-3}}) = 3.$$

In general,

$$\text{tr}(T_{z^m}T_{z^n} - T_{z^n}T_{z^m}) = \begin{cases} n & \text{if } m + n = 0 \\ 0 & \text{otherwise.} \end{cases}$$

Also observe that

$$\frac{1}{2\pi i} \int_0^{2\pi} e^{im\theta} d(e^{in\theta}) = \frac{1}{2\pi i} \int_0^{2\pi} in e^{im\theta} e^{in\theta} d\theta = \begin{cases} n & \text{if } m + n = 0 \\ 0 & \text{otherwise.} \end{cases}$$

**Theorem:** For  $\phi$  and  $\psi$  in  $C^\infty(S^1)$ ,

$$\text{tr}(T_\phi T_\psi - T_\psi T_\phi) = \frac{1}{2\pi i} \int_{S^1} \phi d\psi.$$

**Proof:** Write  $\phi$  and  $\psi$  in terms of the basis  $\{z^n : n \geq 0\}$  and combine the linearity of the trace and the integral with the computations in the example above.  $\square$

We can generalize this result somewhat. Define

$$\mathcal{T}^\infty := \{T_\phi + L : \phi \in C^\infty(S^1), L \in \mathcal{L}^1(H^2(S^1))\}.$$

Then there exists a short exact sequence

$$0 \longrightarrow \mathcal{L}^1(H^2(S^1)) \longrightarrow \mathcal{T}^\infty \xrightarrow{\sigma} C^\infty(S^1) \longrightarrow 0,$$

and the *symbol map*  $\sigma : \mathcal{T}^\infty \rightarrow C^\infty(S^1)$  is given by the formula  $\sigma(T_\phi + L) = \phi$ .

**Theorem:** For  $T_1$  and  $T_2$  in  $\mathcal{T}^\infty$ ,

$$\mathrm{tr}(T_1T_2 - T_2T_1) = \frac{1}{2\pi i} \int_{S^1} \sigma(T_1) d(\sigma(T_2)).$$

**Proof:** Write  $T_1 = T_{\phi_1} + L_1$  and  $T_2 = T_{\phi_2} + L_2$ . Then

$$\begin{aligned} \mathrm{tr}(T_1T_2 - T_2T_1) &= \mathrm{tr}((T_{\phi_1} + L_1)(T_{\phi_2} + L_2) - (T_{\phi_2} + L_2)(T_{\phi_1} + L_1)) \\ &= \mathrm{tr}(T_{\phi_1}T_{\phi_2} - T_{\phi_2}T_{\phi_1} + T_{\phi_1}L_2 - L_2T_{\phi_1} \\ &\quad + L_1T_{\phi_2} - T_{\phi_2}L_1 + L_1L_2 - L_2L_1) \\ &= \mathrm{tr}(T_{\phi_1}T_{\phi_2} - T_{\phi_2}T_{\phi_1}) + \mathrm{tr}(T_{\phi_1}L_2 - L_2T_{\phi_1}) \\ &\quad + \mathrm{tr}(L_1T_{\phi_2} - T_{\phi_2}L_1) + \mathrm{tr}(L_1L_2 - L_2L_1) \\ &= \mathrm{tr}(T_{\phi_1}T_{\phi_2} - T_{\phi_2}T_{\phi_1}) \\ &= \frac{1}{2\pi i} \int \phi_1 d\phi_2 \\ &= \frac{1}{2\pi i} \int \sigma(T_1) d(\sigma(T_2)). \end{aligned}$$

Note that  $\mathrm{tr}(T_1T_2 - T_2T_1)$  only depends on the symbols of  $T_1$  and  $T_2$ !

Now let's look at the determinant.

Suppose  $\phi, \psi$  are nowhere-vanishing functions in  $C^\infty(S^1)$  and that the winding numbers of  $\phi$  and  $\psi$  are zero. Then  $T_\phi$  and  $T_\psi$  are invertible (this is a nontrivial fact!).

Warning:  $T_\phi^{-1} \neq T_{\phi^{-1}}$  in general!

Note that

$$\sigma(T_\phi T_\psi T_\phi^{-1} T_\psi^{-1}) = \phi\psi\phi^{-1}\psi^{-1} = 1,$$

whence  $T_\phi T_\psi T_\phi^{-1} T_\psi^{-1}$  is in  $I + \mathcal{L}^1(H^2(S^1))$ .

$$\det(T_\phi T_\psi T_\phi^{-1} T_\psi^{-1}) = ??$$

It's not too hard to prove that the quantity we are taking the determinant of only depends on the symbols  $\phi$  and  $\psi$ . That is, if  $T_1$  and  $T_2$  are invertible Toeplitz operators with  $\sigma(T_1) = \phi$  and  $\sigma(T_2) = \psi$ , then

$$\det(T_1 T_2 T_1^{-1} T_2^{-1}) = \det(T_\phi T_\psi T_\phi^{-1} T_\psi^{-1}).$$

**Theorem [Campbell-Baker-Hausdorff-Dynkin-...]:** Suppose  $A$  and  $B$  are operators on  $\mathcal{H}$ . If  $\|A\| + \|B\| < \sqrt{2}$ , then  $(\exp A)(\exp B) = \exp C$ , where

$$C = A + B + \frac{1}{2}[A, B] + \frac{1}{12}[A, [A, B]] - \frac{1}{12}[B, [A, B]] \\ + \text{terms involving higher commutators of } A \text{ and } B.$$

**Corollary:**

$$(\exp A)(\exp B)(\exp(-A))(\exp(-B)) = \\ \exp([A, B] + \text{terms involving higher commutators of } A \text{ and } B).$$

Suppose  $\phi$  and  $\psi$  are close to 1. Then  $\phi = e^\alpha$  and  $\psi = e^\beta$  for functions  $\alpha$  and  $\beta$  in  $C^\infty(S^1)$ . Note that

$$\sigma(\exp T_\alpha) = e^{\sigma(T_\alpha)} = e^\alpha = \phi \\ \sigma(\exp T_\beta) = e^{\sigma(T_\beta)} = e^\beta = \psi.$$

Therefore

$$\det(T_\phi T_\psi T_\phi^{-1} T_\psi^{-1}) = \det(\exp(T_\alpha) \exp(T_\beta) \exp(-T_\alpha) \exp(-T_\beta)) \\ = \det \exp(T_\alpha T_\beta - T_\beta T_\alpha + \text{higher commutators}) \\ = \exp \operatorname{tr}(T_\alpha T_\beta - T_\beta T_\alpha + \text{higher commutators}) \\ = \exp\left(\frac{1}{2\pi i} \int_{S^1} \alpha d\beta\right) \\ = \exp\left(\frac{1}{2\pi i} \int_{S^1} \log \phi d(\log \psi)\right) \\ = \exp\left(\frac{1}{2\pi i} \int_{S^1} \log \phi \cdot \frac{d\psi}{\psi}\right).$$

Let's look at this from a different point of view.

Let  $\mathcal{H}$  be a Hilbert space. Then  $\mathcal{H}^n$  is also a Hilbert space:

$$\langle (v_1, v_2, \dots, v_n), (w_1, w_2, \dots, w_n) \rangle := \langle v_1, w_1 \rangle + \langle v_2, w_2 \rangle + \dots + \langle v_n, w_n \rangle.$$

We can view elements of  $\mathcal{B}(\mathcal{H}^n)$  as elements of  $M(n, \mathcal{B}(\mathcal{H}))$ . By extending the notion of symbol in the obvious way, we have a short exact sequence

$$0 \longrightarrow \mathcal{L}^1((H^2(S^1))^n) \longrightarrow M(n, \mathcal{T}^\infty) \xrightarrow{\sigma} M(n, C^\infty(S^1)) \longrightarrow 0.$$

Suppose  $\phi$  and  $\psi$  are arbitrary invertible elements of  $C^\infty(S^1)$ . Then we can find matrices  $R$  and  $S$  in  $GL(3, \mathcal{T}^\infty)$  such that

$$\sigma(R) = \begin{pmatrix} \phi & 0 & 0 \\ 0 & \phi^{-1} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

and

$$\sigma(S) = \begin{pmatrix} \psi & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & \psi^{-1} \end{pmatrix}.$$

For example, we can choose

$$R = \begin{pmatrix} 2T_\phi - T_\phi T_{\phi^{-1}} T_\phi & T_\phi T_{\phi^{-1}} - I & 0 \\ I - T_{\phi^{-1}} T_\phi & T_{\phi^{-1}} & 0 \\ 0 & 0 & I \end{pmatrix}$$

and

$$S = \begin{pmatrix} 2T_\psi - T_\psi T_{\psi^{-1}} T_\psi & 0 & T_\psi T_{\psi^{-1}} - I \\ 0 & I & 0 \\ I - T_{\psi^{-1}} T_\psi & 0 & T_{\psi^{-1}} \end{pmatrix}$$

We infer from the short exact sequence above that the operator  $RSR^{-1}S^{-1}$  is determinant-class. Furthermore, the value of this determinant does not depend on the choice of  $R$  and  $S$  satisfying the properties above - the determinant of  $RSR^{-1}S^{-1}$  only depends on  $\phi$  and  $\psi$ .

Suppose that  $\phi$  and  $\psi$  are restrictions of meromorphic functions (which we also denote  $\phi$  and  $\psi$ ) defined in a neighborhood of the closed unit disk such that neither  $\phi$  nor  $\psi$  has zeros or poles on the unit circle. For each point  $z$  in the open unit disk  $\mathbb{D}$ , define

$$v(\phi, z) = \begin{cases} m & \text{if } \phi \text{ has a zero of order } m \text{ at } z \\ -m & \text{if } \phi \text{ has a pole of order } m \text{ at } z \\ 0 & \text{if } \phi \text{ has neither a zero nor a pole at } z, \end{cases}$$

and similarly define  $v(\psi, z)$ . The quantity

$$\lim_{w \rightarrow z} (-1)^{v(\phi, z)v(\psi, z)} \frac{\psi(w)^{v(\phi, z)}}{\phi(w)^{v(\psi, z)}}$$

is called the *tame symbol* of  $\phi$  and  $\psi$  at  $z$  and is denoted  $(\phi, \psi)_z$ .

Example:

$$\phi(z) = \frac{z^3 - 3z^2}{2z + 1} \quad \text{double zero at 0, simple zero at 3, simple pole at } -1/2$$

$$\psi(z) = \frac{2z - 1}{z^3} \quad \text{simple zero at } 1/2, \text{ triple pole at } 0$$

$$\begin{aligned} (\phi, \psi)_0 &= \lim_{w \rightarrow 0} \left( (-1)^{(2)(-3)} \frac{\left(\frac{2w-1}{w^3}\right)^2}{\left(\frac{w^2(w-3)}{2w+1}\right)^{-3}} \right) \\ &= \lim_{w \rightarrow 0} \frac{(2w-1)^2}{w^6} \cdot \frac{w^6(w-3)^3}{(2w+1)^3} \\ &= \lim_{w \rightarrow 0} \frac{(2w-1)^2(w-3)^3}{(2w+1)^3} \\ &= -27 \end{aligned}$$

$$\begin{aligned} (\phi, \psi)_{-1/2} &= \lim_{w \rightarrow -1/2} \left( (-1)^{(-1)(0)} \frac{\left(\frac{2w-1}{w^3}\right)^{-1}}{\left(\frac{w^2(w-3)}{2w+1}\right)^0} \right) \\ &= \lim_{w \rightarrow -1/2} \frac{w^3}{2w-1} \\ &= \frac{1}{16} \end{aligned}$$

$$\begin{aligned} (\phi, \psi)_{1/2} &= \lim_{w \rightarrow 1/2} \left( (-1)^{(0)(-1)} \frac{\left(\frac{2w-1}{w^3}\right)^0}{\left(\frac{w^2(w-3)}{2w+1}\right)^1} \right) \\ &= \lim_{w \rightarrow 1/2} \frac{2w+1}{w^2(w-3)} \\ &= -\frac{16}{5} \end{aligned}$$

We will not compute  $(\phi, \psi)_3$  for reasons that will become clear in a minute. For all other complex numbers  $z$ , we see that  $(\phi, \psi)_z = 1$ .

**Theorem:**

$$\det(RSR^{-1}S^{-1}) = \prod_{z \in \mathbb{D}} (\phi, \psi)_z^{-1}.$$

Remark 1: Suppose that  $T_\phi$  and  $T_\psi$  are invertible. Then we can take

$$R = \begin{pmatrix} T_\phi & 0 & 0 \\ 0 & T_\phi^{-1} & 0 \\ 0 & 0 & I \end{pmatrix}$$

and

$$S = \begin{pmatrix} T_\psi & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & T_\psi^{-1} \end{pmatrix},$$

whence

$$\det(RSR^{-1}S^{-1}) = \det \begin{pmatrix} T_\phi T_\psi T_\phi^{-1} T_\psi^{-1} & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & I \end{pmatrix} = \det \left( T_\phi T_\psi T_\phi^{-1} T_\psi^{-1} \right).$$

Remark 2: In fact,  $\det(RSR^{-1}S^{-1})$  only depends on the *Steinberg symbol*  $\{\phi, \psi\}$  of  $\phi$  and  $\psi$ . This is an element of the algebraic  $K$ -theory group  $K_2(C^\infty(S^1))$ , and we can use the above theorem to prove that certain Steinberg symbols are nontrivial.

Surprising fact that comes out of this circle of ideas: if both  $\phi$  and  $\psi := 1 - \phi$  are invertible, then  $\det(RSR^{-1}S^{-1}) = 1$ .

The de la Harpe-Skandalis “Determininant”

Suppose  $A$  is a unital Banach algebra with a trace  $\tau : A \rightarrow \mathbb{C}$ . Then we can extend  $\tau$  to a trace on  $M(n, A)$  in the obvious way.

Let  $GL_0(n, A)$  denote the connected component of the identity matrix in  $GL(n, A)$ . Then given a  $C^1$ -path  $\xi$  in  $GL_0(n, A)$ , define

$$\tilde{\Delta}(\xi) = \tau \left( \frac{1}{2\pi i} \int_0^1 \xi'(t) \xi(t)^{-1} dt \right) = \frac{1}{2\pi i} \int_0^1 \tau (\xi'(t) \xi(t)^{-1}) dt.$$

Properties of  $\tilde{\Delta}$ :

- Suppose that  $\xi = \xi_1 \cdot \xi_2$  (pointwise product). Then  $\tilde{\Delta}(\xi) = \tilde{\Delta}(\xi_1) + \tilde{\Delta}(\xi_2)$ ;
- If  $\|\xi(t) - 1\| < 1$  for all  $0 \leq t \leq 1$ , then

$$\tilde{\Delta}(\xi) = \frac{1}{2\pi i} \tau \left( \log(\xi(1)) - \frac{1}{2\pi i} \log \xi(0) \right);$$

- The value of  $\tilde{\Delta}(\xi)$  only depends on the homotopy class of  $\xi$  *with the endpoints fixed*;
- Given an idempotent  $p$  (i.e.  $p^2 = p$ ) in  $M(n, A)$ , define  $\xi_p$  by the formula  $\xi_p(t) = e^{2\pi i t} p + (1 - p)$ . Then  $\tilde{\Delta}(\xi_p) = \tau(p)$ .

Suppose  $x$  is an element of  $GL_0(n, A)$ , choose a  $C^1$ -path  $\xi$  from 1 to  $x$ , and define  $\Delta(x) = \tilde{\Delta}(\xi)$ .

Problem:  $\Delta(x)$  depends on the choice of path  $\xi$ .

What we really have, via the (in)famous Bott periodicity theorem in  $K$ -theory, is a function into  $\mathbb{C}/\tau(K_0(A))$ .

Properties of  $\Delta$ :

- $\Delta$  is a group homomorphism;
- $\Delta$  is surjective if and only if  $\tau$  is surjective;
- $\Delta(e^y) = \tau(y) + \tau(K_0(A))$  for  $y$  in  $M(n, A)$ .

**Corollary:** Suppose that  $\tau(K_0(A)) \cong \mathbb{Z}$ . Then

$$\exp(2\pi i\Delta) : \mathrm{GL}_0(n, A) \rightarrow \mathbb{C}^*$$

is a group homomorphism, and

$$\exp(2\pi i\Delta)(e^y) = e^{\tau(y)}$$

for  $y$  in  $M(n, A)$ . In particular, if  $A = \mathbb{C}$  and  $\tau$  is the identity map, then  $\exp(2\pi i\Delta)$  is the usual determinant on  $M(n, \mathbb{C})$ .

The Fuglede-Kadison-Brown-Hochs-Kaad-Schemaitat Determinant

A *von Neumann algebra* is a  $C^*$ -subalgebra of  $\mathcal{B}(\mathcal{H})$  that is closed in the topology of pointwise convergence.

Example:  $L^\infty(\mathbb{R})$

Suppose we have a “nice” (normal, faithful, semifinite) trace  $\tau$  defined on positive elements in  $\mathcal{N}$ .

Example: For  $L^\infty(\mathbb{R})$ , take  $\tau(f) = \int_{-\infty}^{\infty} f(x) dx$ .

Let  $\mathcal{L}^1(\mathcal{N})$  be the trace ideal. For invertible elements in  $\mathcal{N}$  of the form  $1+x$  with  $x$  in  $\mathcal{L}^1(\mathcal{N})$ , the aforementioned authors define a determinant homomorphism  $\det_\tau$  with values in  $(0, \infty)$ :

$$\det_\tau(1+x) = e^{\tau(\log |1+x|)}.$$

This determinant is multiplicative – this is highly nontrivial to show and involves techniques from algebraic  $K$ -theory and Connes’ cyclic homology.



Example: Wiener-Hopf Operators

Consider the Hilbert space  $L^2(\mathbb{R})$ . There is a Hilbert subspace  $H^2(\mathbb{R})$  of  $L^2(\mathbb{R})$  that consists of elements that have an analytic extension to the upper half plane, satisfying a certain growth condition. Let  $P$  denote orthogonal projection from  $L^2(\mathbb{R})$  onto  $H^2(\mathbb{R})$ .

Let  $C_b(\mathbb{R})$  denote the algebra of bounded continuous functions on the real line. For each  $\phi$  in  $C_b(\mathbb{R})$ , multiplication by  $\phi$  defines an element of  $\mathcal{B}(L^2(\mathbb{R}))$ , and we can compress to  $H^2(\mathbb{R})$  just as we did in the circle case to obtain an operator  $W_\phi$  on  $\mathcal{B}(H^2(\mathbb{R}))$ :

$$W_\phi = PM_\phi.$$

The algebra of *almost periodic* functions on  $\mathbb{R}$  is the norm-closed subalgebra of  $C_b(\mathbb{R})$  generated by the functions  $t \rightarrow e^{i\lambda t}$  for  $\lambda$  real. This algebra can be identified with  $C(\mathbb{R}_B)$ , the continuous functions on the *Bohr compactification* of  $\mathbb{R}$ .

Let  $\mathcal{W}$  be the  $C^*$ -subalgebra of  $\mathcal{B}(H^2(\mathbb{R}))$  generated  $\{W_\phi : \phi \in C(\mathbb{R}_B)\}$ . Then there is a short exact sequence

$$0 \longrightarrow \mathcal{C} \longrightarrow \mathcal{W} \xrightarrow{\sigma} C(\mathbb{R}_B) \longrightarrow 0,$$

where  $\mathcal{C}$  is the commutator ideal of  $\mathcal{W}$ . This commutator ideal lives inside, but is not equal to, the trace ideal associated to a von Neumann trace. Just as in the circle case, there is a “smooth” version of this short exact sequence:

$$0 \longrightarrow \mathcal{C}^\infty \longrightarrow \mathcal{W}^\infty \xrightarrow{\sigma} C^\infty(\mathbb{R}_B) \longrightarrow 0.$$

**Theorem:** Suppose  $W_1$  and  $W_2$  in  $\mathcal{W}^\infty$  have symbols  $\phi$  and  $\psi$  respectively. Then

$$\mathrm{tr}(W_1W_2 - W_2W_1) = \lim_{R \rightarrow \infty} \left( \frac{1}{2R} \int_{-R}^R \phi(t)\psi'(t) dt \right).$$

**Theorem:** Suppose  $W_1$  and  $W_2$  in  $\mathcal{W}^\infty$  are invertible, have symbols  $\phi$  and  $\psi$  respectively, and are close to  $I$ . Then

$$\det_\tau(W_1W_2W_1^{-1}W_2^{-1}) = \lim_{R \rightarrow \infty} \left( \exp \left( \frac{1}{2R} \mathrm{Re} \left( \int_{-R}^R \log(\phi(t)) \frac{\psi'(t)}{\psi(t)} dt \right) \right) \right).$$