

Speculations on Mirror Symmetry and the Arithmetic of the Quintic Threefold

with

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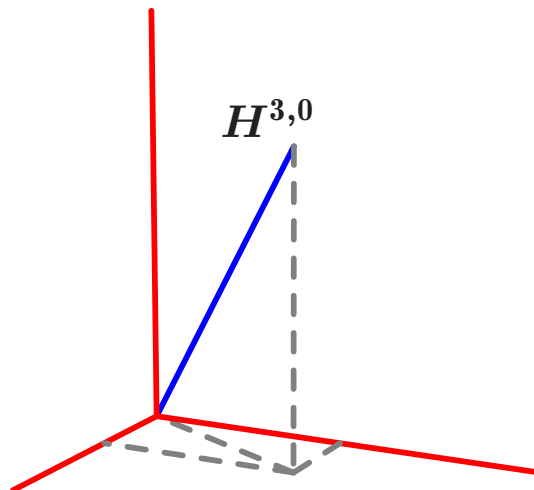
Calabi-Yau manifolds arise when we ‘compactify’ a $D = 10$ string theory down to $D = 4$. These manifolds have many remarkable properties owing to this relation with string theory and supersymmetry.

- The number of \mathbb{F}_p -rational points may be calculated in terms of periods. Moreover this expression is computable in a practical sense.
- We are led to compute the corresponding ζ -function and to speculate on the possible role of ‘quantum corrections’ in relation to mirror symmetry and the conifold singularization.

An important property of a Calabi–Yau manifold is that there exists a unique holomorphic three-form, Ω . One way to parametrise the complex structure is via the **algebraic parameters** as in the quintic threefold

$$P = x_1^5 + x_2^5 + x_3^5 + x_4^5 + x_5^5 + \alpha x_1^2 x_2^3 + \beta x_4^3 x_3 x_5 + \dots .$$

For the quintic there are $h^{2,1} = 101$ such parameters.



The (projective) coordinates for $H^{3,0}$ are the **periods**

$$\varpi_j(\psi) = \int_{\gamma_j} \Omega .$$

Consider, for definiteness, the one parameter family of quintics in \mathbb{P}_4

$$\mathcal{M} : P(x, \psi) = \sum_{i=1}^5 x_i^5 - 5\psi x_1 x_2 x_3 x_4 x_5 .$$

\mathcal{M} has $h^{11} = 1$ and $h^{21} = 101$. It also has a group $\Gamma \cong \mathbb{Z}_5^3$ of automorphisms:

$$(x_1, x_2, x_3, x_4, x_5) \mapsto (\zeta^{n_1} x_1, \zeta^{n_2} x_2, \zeta^{n_3} x_3, \zeta^{n_4} x_4, \zeta^{n_5} x_5)$$

with $\zeta^5 = 1$ and $\sum n_i = 0$.

In this simple case there is a simple relation between \mathcal{M} and its mirror

$$\mathcal{W} = \mathcal{M}/\Gamma .$$

\mathcal{W} has $h^{11} = 101$ and $h^{21} = 1$. The 100 extra Kähler forms arise from the blow ups required to resolve the singularities associated with the fixed points of Γ .

Now, as we have said, the number of rational points is determined by the periods and there are $b^3 = 2h^{21} + 2$ of these. The Hodge number h^{21} counts the number of parameters on which the complex structure depends and, in simple cases, this corresponds to the number of ways of deforming the defining polynomial

$$P(x, c) = \sum_v c_v x^v \quad ; \quad x^v = x_1^{v_1} x_2^{v_2} x_3^{v_3} x_4^{v_4} x_5^{v_5} .$$

The directions in which $P(x, c)$ can be deformed correspond to the monomials x^v considered subject to the ideal $(\partial P / \partial x_i)$. A special role is played by fundamental monomial

$$Q = x_1 x_2 x_3 x_4 x_5$$

which is related by mirror symmetry to the Kähler form of the mirror.

Return now to our special one parameter family of polynomials

$$P(x, \psi) = \sum_{i=1}^5 x_i^5 - 5\psi x_1 x_2 x_3 x_4 x_5 .$$

\mathcal{M} has $2h^{21}(\mathcal{M}) + 2 = 204 = 2 \times 100 + 4$ periods while \mathcal{W} has $2h^{21}(\mathcal{W}) + 2 = 4$.

$$\begin{array}{ccccccc} 1 & \longrightarrow & Q & \longrightarrow & Q^2 & \longrightarrow & Q^3 \\ & & x^v & \longrightarrow & Q x^v & & \end{array}$$

This leads to 1 fourth order differential operator \mathcal{L}_0 and 100 second order operators \mathcal{L}_v .

There are tenth order monomials that are not included in the above scheme and which require special attention. The generators of the ideal are

$$x_1^4 \simeq \psi x_2 x_3 x_4 x_5 \text{ \& cyclic.}$$

Thus

$$x^{(4,3,2,1,0)} \simeq \psi x^{(0,4,3,2,1)} \simeq \dots \simeq \psi^5 x^{(4,3,2,1,0)} .$$

The Z-Function

We now work over \mathbb{F}_{p^r} and let $N_r(\psi)$ denote the number of projective solutions to $P(x, \psi) = 0$. Our expressions lead us to decompose N_r into a sum of contributions $N_r = N_{r,0} + \sum_v N_{r,v}$.

The ζ -function is defined by the expression

$$\zeta(t) = \exp \left(\sum_{r=1}^{\infty} \frac{N_r t^r}{r} \right)$$

Given the decomposition of N_r into a sum over monomials v we have

$$\zeta_{\mathcal{M}}(t, \psi) = \frac{R_0(t, \psi) \prod_v R_v(t, \psi)}{(1-t)(1-pt)(1-p^2t)(1-p^3t)}$$

$$\zeta_{\mathcal{W}}(t, \psi) = \frac{R_0(t, \psi)}{(1-t)(1-pt)^{101}(1-p^2t)^{101}(1-p^3t)} .$$

In all cases, apart from the conifold, R_0 is a quartic

$$R_0 = 1 + a_0 t + b_0 p t^2 + a_0 p^3 t^3 + p^6 t^4 .$$

The Euler Curves

Classical analysis gives an expression for the hypergeometric functions in terms of Euler's integral which is of the form

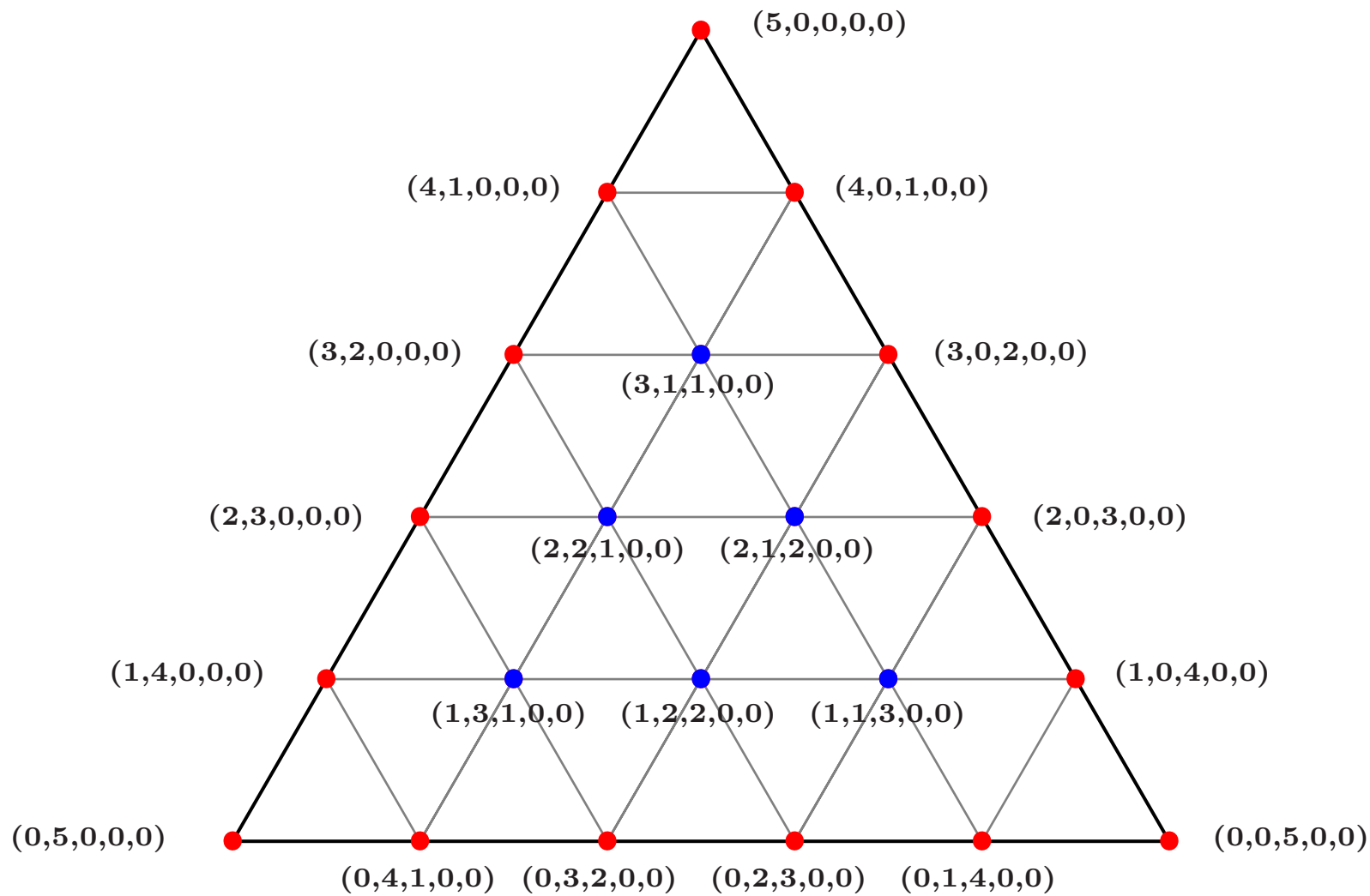
$$\int dx x^{-\alpha/5} (1-x)^{-\beta/5} (1-x/\psi^5)^{-(1-\beta/5)} .$$

If we think of Euler's integral as $\int \frac{dx}{y}$ then we are led to curves

$$\mathcal{E}_{\alpha\beta}(\psi) : y^5 = x^\alpha (1-x)^\beta (1-x/\psi^5)^{5-\beta} .$$

v	α	β
(4, 1, 0, 0, 0)	2	3
(3, 2, 0, 0, 0)	1	4
(3, 1, 1, 0, 0)	2	4
(2, 2, 1, 0, 0)	4	3

$$\mathcal{E}_{\alpha\beta} = \begin{cases} \mathcal{A} & \alpha + \beta = 5 \\ \mathcal{B} & \alpha + \beta \neq 5 \text{ and } \alpha \neq \beta . \end{cases}$$



For the curve \mathcal{A} there is a corresponding ζ -function

$$\zeta_{\mathcal{A}}(t) = \frac{R_{\mathcal{A}}(t)^2}{(1-t)(1-pt)} .$$

Now the existence of nontrivial fifth roots of unity is important for the mirror construction. Such roots of unity exist in \mathbb{F}_{p^r} precisely when $5|p^r - 1$. For given p let $\rho = 1, 2$ or 4 be the smallest r for which $5|p^r - 1$.

The R_v pair up in the following way:

$$R_{(4,1,0,0,0)}(t) R_{(3,2,0,0,0)}(t) = R_{\mathcal{A}}(p^\rho t^\rho)^{1/\rho}$$

$$R_{(3,1,1,0,0)}(t) R_{(2,2,1,0,0)}(t) = R_{\mathcal{B}}(p^\rho t^\rho)^{1/\rho} .$$

So the ζ -function for \mathcal{M} takes the form

$$\zeta_{\mathcal{M}}(t, \psi) = \frac{R_0(t, \psi) R_{\mathcal{A}}(p^\rho t^\rho, \psi)^{\frac{30}{\rho}} R_{\mathcal{B}}(p^\rho t^\rho, \psi)^{\frac{20}{\rho}}}{(1-t)(1-pt)(1-p^2t)(1-p^3t)} .$$

The Z-Function and Mirror Symmetry

As defined the ζ -function does not respect mirror symmetry

$$\zeta(t) = \frac{\text{Numerator of deg. } 2h^{21} + 2 \text{ depending on the cpx. structure of } \mathcal{M}}{\text{Denominator of deg. } 2h^{11} + 2}.$$

Explicitly for the quintic we have

$$\zeta_{\mathcal{M}}(t, \psi) = \frac{R_0(t, \psi) R_{\mathcal{A}}(pt, \psi)^{30} R_{\mathcal{B}}(pt, \psi)^{20}}{(1-t)(1-pt)(1-p^2t)(1-p^3t)}$$

$$\zeta_{\mathcal{W}}(t, \psi) = \frac{R_0(t, \psi)}{(1-t)(1-pt)^{101}(1-p^2t)^{101}(1-p^3t)}$$

Clearly what we would like to see is

$$\zeta_{\mathcal{W}} = \frac{1}{\zeta_{\mathcal{M}}}$$

however we cannot have this relation and the classical definition in terms of the N_r with positive N_r .

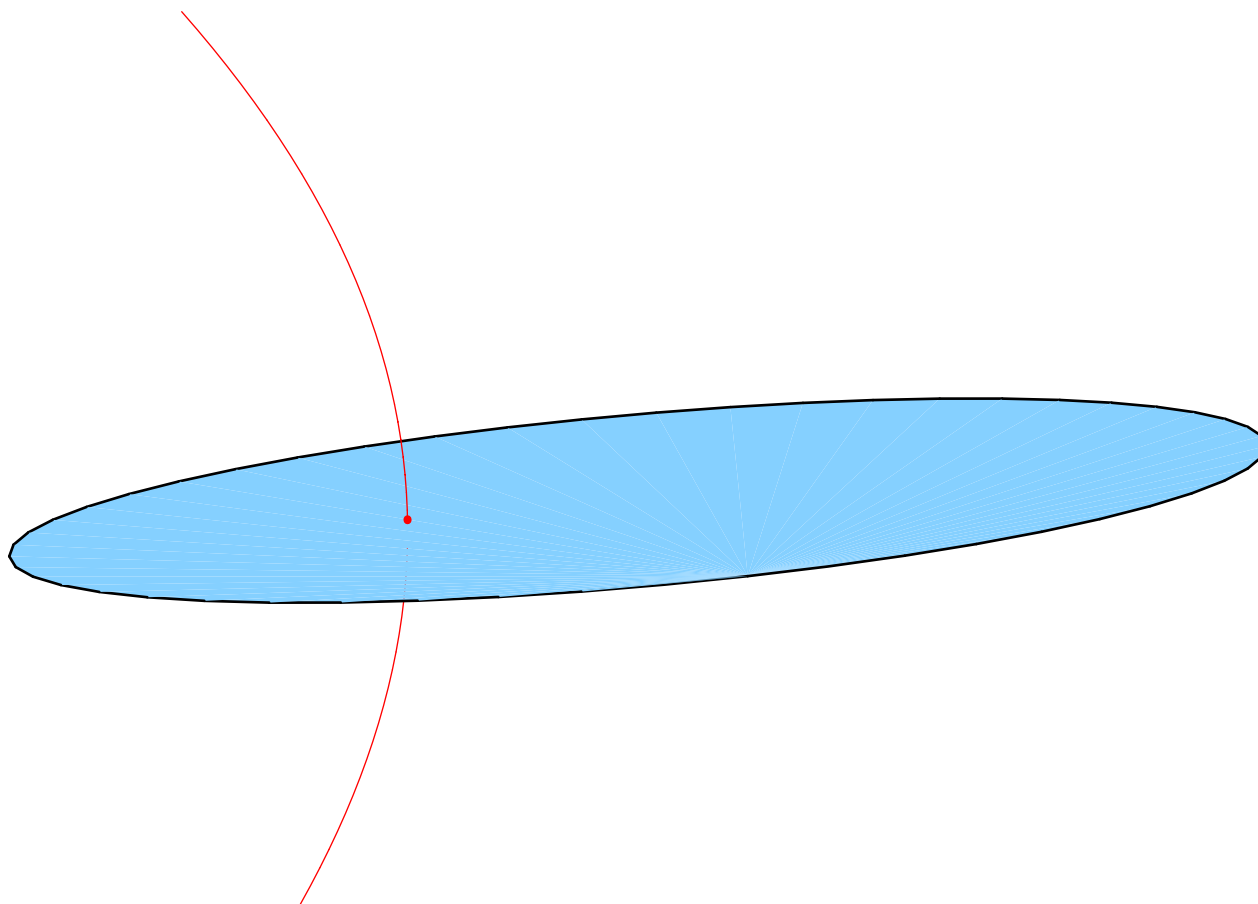
The Conifold

For the conifold $\psi^5 = 1$ the ζ -function seems to be especially simple

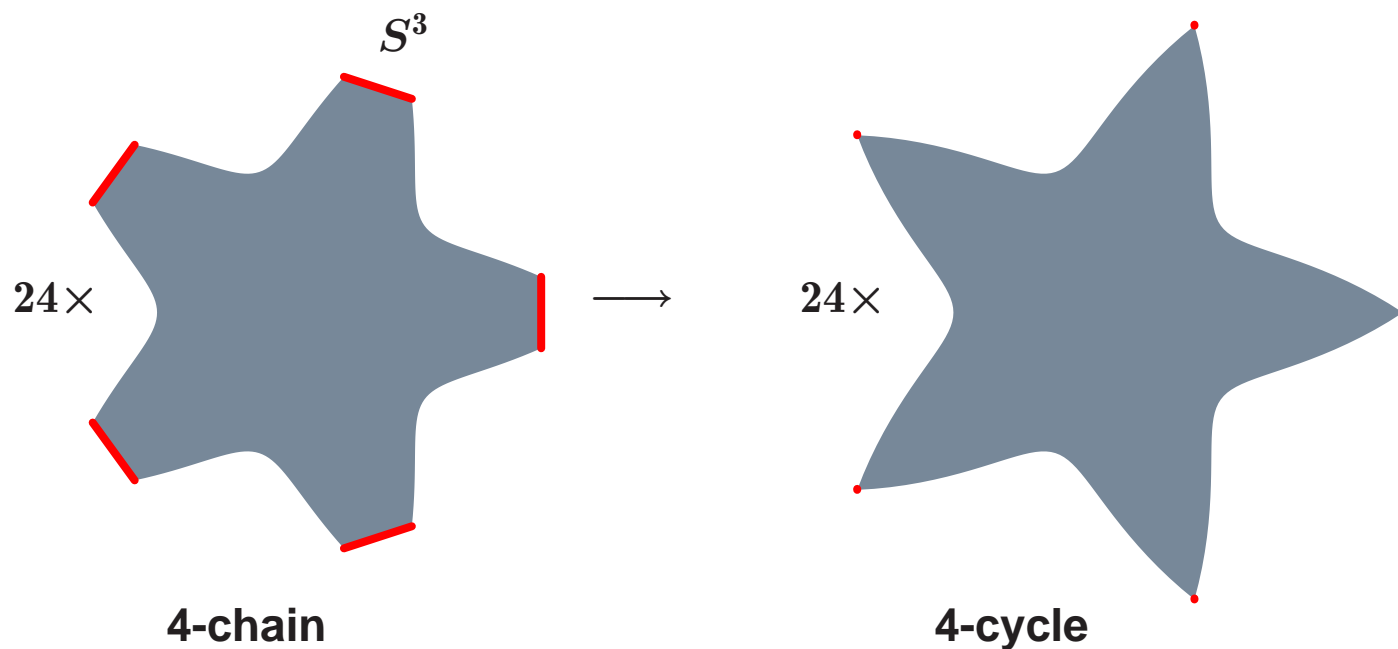
$$\zeta(t, \psi^5 = 1) = \frac{(1 - \epsilon pt) (1 - a_p t + p^3 t^2) (1 - p^\rho t^\rho)^{\frac{200}{\rho}}}{(1 - t)(1 - pt)(1 - p^2 t)(1 - p^3 t) (1 - p^{2\rho} t^\rho)^{\frac{24}{\rho}}}$$

where $\epsilon = \left(\frac{5}{p}\right) = \pm 1$ and a_p is the p -th coefficient in the q -expansion of the eigenform, g , found by Schoen; it is the unique cusp form of weight 4 for the group $\Gamma_0(25)$.

$$\begin{aligned} g &= \eta(q^5)^4 [\eta(q)^4 + 5 \eta(q)^3 \eta(q^{25}) + 20 \eta(q)^2 \eta(q^{25})^2 + 25 \eta(q) \eta(q^{25})^3 + 25 \eta(q^{25})^4] \\ &= q + q^2 + 7 q^3 - 7 q^4 + 7 q^6 + 6 q^7 - 15 q^8 + 22 q^9 - 43 q^{11} - 49 q^{12} - 28 q^{13} \\ &\quad + 6 q^{14} + 41 q^{16} + 91 q^{17} + 22 q^{18} - 35 q^{19} + 42 q^{21} - 43 q^{22} + 162 q^{23} \\ &\quad - 105 q^{24} - 28 q^{26} - 35 q^{27} - 42 q^{28} + 160 q^{29} + 42 q^{31} + \dots \end{aligned}$$

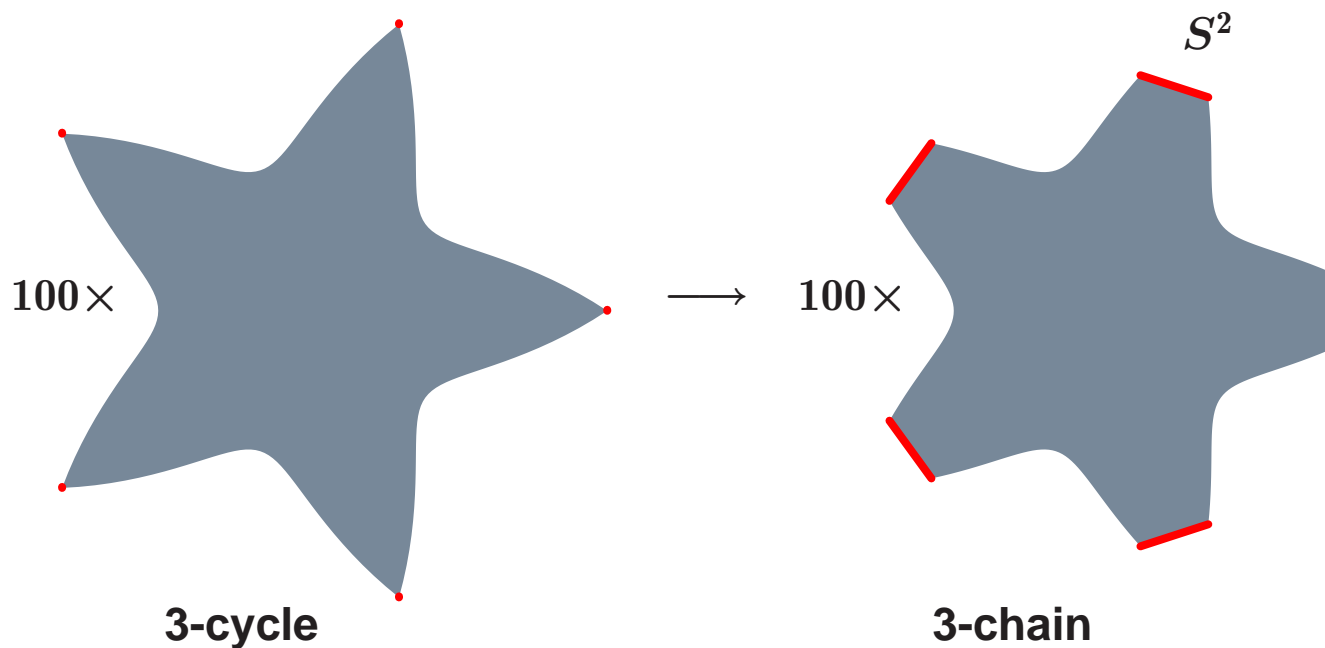


125 S^3 's are blown down but only 101 are independent so 24 4-cycles are created.



$$\zeta(t, \psi^5 = 1) = \frac{(1 - a_p t + p^3 t^2) (1 - pt)^{200}}{(1 - t)(1 - p^2 t)^{25}(1 - p^3 t)}$$

Now we resolve 125 nodes with \mathbb{P}^1 's, but there are 100 relations so we destroy 100 3-cycles.



$$\begin{aligned} \zeta(t, \psi^5 = 1) &= \frac{(1 - a_p t + p^3 t^2) (1 - pt)^{200-100}}{(1 - t)(1 - pt)^{125} (1 - p^2 t)^{25} (1 - p^3 t)} \\ &= \frac{(1 - a_p t + p^3 t^2)}{(1 - t)(1 - pt)^{25} (1 - p^2 t)^{25} (1 - p^3 t)}. \end{aligned}$$

